## Protein dynamics and markov modeling



#### Frank Noé

Talk 01 - Introduction + Overview





Before we start...

# installing for the first time? conda config --add channels conda-forge

# install / upgrade PyEMMA conda install pyemma

test your installation: import pyemma print pyemma.\_\_version\_\_











# Proteins





McGufee and Elcock, PloS Comput Biol 2010





# Protein-Protein binding









# Protein-Protein binding





Plattner, Doerr, De Fabritiis, Noé

0.1 microseconds





# 50 K atom system (all atom, explicit solvent)



Rate

350 ns / day / GPU\* e.g. Amber, AceMD, OpenMM on Titan X 70 µs / day / Anton II





# 50 K atom system (all atom, explicit solvent)



| Rate       | 350 ns / day / GPU*<br>e.g. Amber, AceMD, OpenMM on Titan X | 70 µs / day / Anton I  |  |
|------------|---|------------------------|--|
|            | 200 GPUs  | 1 Anton II             |  |
|            | 100 traj. of 350 ns / day                                   | 1 traj. of 10 µs / day |  |
| Throughput | 70 µs / day   | 70 μs / day            |  |





# 50 K atom system (all atom, explicit solvent)



350 ns / day / GPU\* e.g. Amber, AceMD, OpenMM on Titan X 70 µs / day / Anton II

|            | 200 GPUs                  | 1 Anton II             |  |
|------------|---------------------------|------------------------|--|
|            | 100 traj. of 350 ns / day | 1 traj. of 10 µs / day |  |
| Throughput | 70 µs / day               | 70 µs / day            |  |
| Cost       | 200.000 USD               | 20.000.000 USD ???     |  |





# Conformation Dynamics / Markov models















So what do we do?











- Molecular motion is primarily driven by thermal fluctuations, and thus inherently stochastic
- A molecular system driven by thermal motion is also reversible and stationary, at least in between "nonreversible checkpoints".
- The stationary distribution is given by the Boltzmann distribution

$$\mu(\mathbf{x}) = Z^{-1} \exp\left(-\frac{U(\mathbf{x})}{k_B T}\right)$$
$$Z = \int_{\mathbf{x} \in \Omega} \exp\left(-\frac{U(\mathbf{x})}{k_B T}\right) \, \mathrm{d}\mathbf{x}$$

Z or integrals over sets of x cannot be computed exactly for nontrival systems, and must therefore be sampled.





Meaningful are expectation values:

$$\begin{split} \mathbb{E}[a] &= \int_{\mathbf{x}\in\Omega} \mu(\mathbf{x}) \ a(\mathbf{x}) \ \mathrm{d}\mathbf{x} \\ \mathbb{E}[(a,b);\tau] &= \int_{\mathbf{x}\in\Omega} \int_{\mathbf{y}\in\Omega} \mu(\mathbf{x}) \ a(\mathbf{x}) \ p(\mathbf{x}\to\mathbf{y};\tau) \ b(\mathbf{x}) \ \mathrm{d}\mathbf{x} \ \mathrm{d}\mathbf{y} \end{split}$$

• Example 1: probability of being in the folded state F (rather than unfolded U):

$$p_F = \mathbb{E}[\mathbf{1}_F] = \int_{\mathbf{x}\in\mathbf{F}} \mu(\mathbf{x}) \,\mathrm{d}\mathbf{x}$$

and the free energy difference of folding is then

$$\frac{\Delta G}{k_B T} = -\ln \frac{p_F}{1 - p_F}$$

• Problem: In order to evaluate the above integrals, the parts of state space with significant weights  $\mu(\mathbf{x})$  must be sampled. However, this is very hard because of free energy barriers / metastable states.





# The Markov model trick

We rewrite the problem by introducing a state space partition  $\{S_1, ..., S_n\}$  with  $\Omega = \bigcup_i S_i$ :

$$\mathbb{E}[a] = \sum_{i} \pi_{i} \int_{\mathbf{x} \in S_{i}} \frac{\mu(\mathbf{x})}{\pi_{i}} a(\mathbf{x}) \, \mathrm{d}\mathbf{x} = \sum_{i} \pi_{i} \bar{a}_{i}$$
(1)  
$$\pi_{i} = \int_{\mathbf{x} \in S_{i}} \mu(\mathbf{x}) \, \mathrm{d}\mathbf{x}$$

The first equation has become much easier - the local distribution  $\mu(\mathbf{x})/\pi_i$  is easy to sample if the discrete states  $S_i$  do not contain internal barriers. However the second equation is still as hard. But we can rewrite it as follows:

$$\boldsymbol{\pi}^{\top} = \boldsymbol{\pi}^{\top} \mathbf{P}(\tau) \tag{2}$$

with the transition matrix

$$p_{ij}(\tau) = \int_{\mathbf{x}\in S_i} \int_{\mathbf{y}\in S_j} \frac{\mu(\mathbf{x})}{\pi_i} p(\mathbf{x}\to\mathbf{y};\tau) \,\mathrm{d}\mathbf{x} \,\mathrm{d}\mathbf{y}$$

This is again relatively easy - we need to prepare starting points  $\mathbf{x}$  according to the local distribution  $\mu(\mathbf{x})/\pi_i$ , then simulate for a (usually short) time  $\tau$  and count the transition if it ends of in  $S_j$ .  $p_{ij}(\tau)$  is just the fraction of transitions ending up in  $S_j$  after time  $\tau$  given that we start from  $S_i$ . So we can estimate it without knowing  $\pi_i$ .

We can then reconstruct the unbiased  $\pi = [\pi_i]$  using Eq. (2), and use that in Eq. (1) to compute the expectation value. We have reduce the global sampling problem to a local sampling problem, which is much easier, and we have gained a perfect parallelization of our problem!











# Conformation Dynamics / Markov models



see also works by: Andersen, Caflisch, Chodera, Deuflhard, Dill, Hummer, Pande, Schütte, Stock, Huisinga, Rao, Roux, Levy





## Journal of Computational Physics

Volume 151, Issue 1, 1 May 1999, Pages 146-168



Regular Article

## A Direct Approach to Conformational Dynamics Based on Hybrid Monte Carlo \*

Ch Schüttea, b, A Fischera, W Huisingaa, P Deuflharda, b



# Generation 1 : focus on metastable states



Hierarchical analysis of conformational dynamics in biomolecules: Transition networks of metastable states Frank Noé<sup>1</sup>, Illia Horenko<sup>2</sup>, Christof Schütte<sup>2</sup> and Jeremy C. Smith<sup>3</sup>

+VIEW AFFILIATIONS

J. Chem. Phys. 126, 155102 (2007); http://dx.doi.org/10.1063/1.2714539

# Automatic discovery of metastable states for the construction of Markov models of macromolecular conformational dynamics

John D. Chodera<sup>1</sup>, Nina Singhal<sup>2</sup>, Vijay S. Pande<sup>3</sup>, Ken A. Dill<sup>4</sup> and William C. Swope<sup>5,a)</sup>

+ VIEW AFFILIATIONS

a) Author to whom correspondence should be addressed. Electronic mail: swope@us.ibm.com

J. Chem. Phys. 126, 155101 (2007); http://dx.doi.org/10.1063/1.2714538



# Generation 2: understanding spectral properties of MSMs

Propagator

timescales

1.0

0.8

0.6

0.4

0.2

0.0

1

Eigenvalue  $\lambda_i$ 

(e)

$$p_{\tau}(\mathbf{z}_{\tau}) = \mathcal{P}(\tau) p_0(\mathbf{z}_0)$$

$$p_{\tau}(\mathbf{z}_0, \mathbf{z}_{\tau}) = \mu(\mathbf{z}_{\tau}) + \sum_{i=2}^{\infty} e^{-\kappa_i \tau} \frac{\phi_i(\mathbf{z}_0)}{\mu(\mathbf{z}_0)} \phi_i(\mathbf{z}_{\tau})$$

processes:

Spectral decomposition





# Generation 2: focus on discretizing transfer operator



- \* No systematic error in the equilibrium distribution
- \* Systematic (discretization) error of MSM kinetics depends on eigenfunction approximation quality and lagtime.
- \* Timescales are always underestimated

Sarich, Noé, Schütte: On the approximation quality of Markov state models **Multiscale Model. Simul.** (2010)

Prinz et al.: Markov models of molecular kinetics: generation and validation. J. Chem. Phys. 134, p174105 (2011)



THE JOURNAL OF CHEMICAL PHYSICS 139, 184114 (2013)

#### Projected and hidden Markov models for calculating kinetics and metastable states of complex molecules

Frank Noé,<sup>a)</sup> Hao Wu,<sup>b)</sup> Jan-Hendrik Prinz,<sup>b)</sup> and Nuria Plattner Department of Mathematics and Computer Science, FU Berlin, Arnimallee 6, 14159 Berlin, Germany





# Generation 3: newer developments - VAMPnets



Altmetric: 13 Citations: 1

More detail >>

Article OPEN

# VAMPnets for deep learning of molecular kinetics

Andreas Mardt, Luca Pasquali, Hao Wu 🌡 Frank Noé 🏁

Nature Communications 9, Article number: 5 (2018) doi:10.1038/s41467-017-02388-1 Received: 14 July 2017 Accepted: 22 November 2017 Published online: 02 January 2018





Max. transition probability: 41% Min. transition probability: 0.5%





# **Optimal reaction coordinates?**

Backward propagator

$$\rho_{\tau} = \mathcal{T}(\tau)\rho_0$$

Spectral decomposition  

$$\rho_{\tau} = \sum_{i=1}^{\infty} e^{-\tau \kappa_i} \langle \psi_i \mid \rho_0 \rangle \psi_i$$

10101

Processes:



Nüske et al, **JCTC** 2014



# How to find the slow coordinates?



#### Variational approach of conformation dynamics (VAC)

Noé and Nüske, **Multiscale Model. Simul.** 11, 635-655 (2013) / ArXiv (2012) Nüske et al, **JCTC** 2014

#### Time-lagged independent component analysis (TICA)

Molgedey and Schuster, **PRL** 1994 Perez-Hernandez et al, **JCP**, 139, 1502 (2013)

Schwantes and Pande, JCTC 2013

www.pyemma.org





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#### PCA

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#### Variational Approach



Variational Approach

Noé and Nüske, **MMS** 11, 635-655 (2013) Nüske et al, **JCTC** 10, 1739-1752 (2014) Perez-Hernandez et al, **JCP**, 139, 1502 (2013) Identification of slow molecular order parameters for Markov model construction

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#### PCA

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#### Variational Approach



#### Variational Approach

Noé and Nüske, **MMS** 11, 635-655 (2013) Nüske et al, **JCTC** 10, 1739-1752 (2014)



#### PCA

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#### Variational Approach

Noé and Nüske, MMS 11, 635-655 (2013) Nüske et al, JCTC 10, 1739-1752 (2014)

#### Kinetic map:

Noé and Clementi, **JCTC** 11, 5002-5011 (2015)



1FME peptide - Simulation data from DESRES, Lindorff-Larsen et al, Science 2011





#### **Estimation of transition matrix**

 $T_{ij}(\tau) = \frac{\mathbb{E}[\chi_i(\mathbf{x}(t)) \ \chi_j(\mathbf{x}(t+\tau))]}{\mathbb{E}[\chi_i(\mathbf{x}(t))]} = \frac{c_{ij}^{\text{corr}}(\tau)}{\pi_i},$ 

**Statistical Error** 

$$p(Y|T) = \prod_{k=1}^{n-1} T_{y_k, y_{k+1}} = p(C|T) = \prod_{i,j=1}^m T_{ij}^{c_{ij}}$$

Estimation:

Prinz et al.: **J. Chem Phys.** 134, 174105 (2011) Bowman et al.: **J. Chem Phys.** 131, 124101 (2009) Noé, **J Chem Phys** 128, 244103 (2008)

Linear Error Perturbation: Sinhal, Pande, JCP 2006 Prinz, Smith, Noé, **Multiscale Model. Simul** 2011

Monte Carlo Noé, **J Chem Phys** 128, 244103 (2008) Chodera, Noé, **J Chem Phys** (2010)





## Step 3: Analysis

#### Transition path theory

Stationary probability

$$\boldsymbol{\pi}^T = \boldsymbol{\pi}^T \mathbf{T}(\boldsymbol{\tau}).$$

Committor

$$-q_i^+ + \sum_{k \in I} T_{ik} q_k^+ = -\sum_{k \in B} T_{ik}.$$



Flux

$$f_{ij} = \pi_i q_i^- T_{ij} q_j^+.$$
  
$$f_{ij}^+ = \max\{0, f_{ij} - f_{ji}\}.$$

Metzner, Vanden-Eijnden, Schütte, **MMS** (2009) Noé et al, **PNAS** (2009) Bereszkovskii, Hummer, Szabo, **JCP** (2009)

#### Metastable states (PCCA)

Deuflhard, Weber.: Linear Alg. Appl. 398C, 161 (2005)

#### **Experimental observables**

Noé et al, **PNAS** 108, p 4822 (2011) Lindner et al, **JCP** 139, 175102 (2013)



# Step 4: Coarse-graining



#### Scherer et al. **JCTC** 11, 5525–5542 (2015).



# Markov State Models

# **Review book**

Advances in Experimental Medicine and Biology 797

Gregory R. Bowman Vijay S. Pande Frank Noé *Editors* 

An Introduction to Markov State Models and Their Application to Long Timescale Molecular Simulation

🕘 Springer



# PyEMMA software





# code: www.github.com/markovmodel

# docs: www.pyemma.org

M. K. Scherer, B. Trendelkamp-Schroer, F. Paul, G. Pérez-Hernández, M. Hoffmann, N. Plattner, C. Wehmeyer, J.-H. Prinz, and F. Noé, "PyEMMA 2: A software package for estimation, validation, and analysis of Markov models," **J. Chem. Theory Comput.** 11, 5525–5542 (2015)



# PyEMMA github site

| 🛈 🚔 GitHub, Inc. (US) — https://githut       | b.com/markovmodel/pyemma                                       | ···· ◙ ☆ 🔍 vampnets →                                   |
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## Application to protein-protein association





# Protein-Protein binding









# **Protein-Protein binding**





Plattner, Doerr, De Fabritiis, Noé **Nature Chemistry** (2017) 0.1 microseconds





#### 1) Adaptive molecular dynamics



2) Dimension reduction (10000 => 10) using variational approach

## 3) Discretization using k-means

4) Hidden Markov model based on microstates

Noé et al, JCP 139, 184114 (2013)



# Validation of the model

 crystal structure 1BRS predicted by the most stable HMM state (95% population)

average heavy-atom RMSD 2.1 A



|                                       | Model                                       | 95% confidence interval                                      | Experiment  |
|---------------------------------------|---|--|---|
| Binding free energy                   | 14.8 kcal / mol                             | (12.3 19.3)  | 16.8 kcal/mol   |
| <ul> <li>Association rate</li> </ul>  | <b>0.74 10<sup>8</sup></b> s⁻¹M⁻¹           | (0.72 0.75)  | <b>1·10<sup>8</sup></b> s⁻¹M⁻¹  |
| <ul> <li>Dissociation rate</li> </ul> | <b>2.7 10-</b> <sup>3</sup> s <sup>-1</sup> | (2.8·10 <sup>-6</sup> 1.8·10 <sup>-1</sup> s <sup>-1</sup> ) | ( <b>4.8·10</b> <sup>-5</sup> S <sup>−1</sup> <b>5.0·10</b> <sup>-4</sup> S <sup>−1</sup> ) |
|                                       |   | Pla  | ttner. Doerr. De Fabritiis. Noé   |



## Mutants by first-order perturbation theory





## Estimation (Reversible Markov state model)



Plattner, Doerr, De Fabritiis, Noé **Nature Chemistry** (2017)



















# Protein-Protein binding







# Acknowledgements



#### Collaborations

Cecilia Clementi (Rice University) Christof Schütte (FU Berlin) Eric Vanden-Eijnden (Courant NY) Thomas Weikl (MPI Potsdam) Edina Rosta (King's College London) Bettina Keller (FU Berlin)

Vijay Pande (Stanford) Volker Haucke (FMP Berlin) Stephan Sigrist (FU Berlin) Oliver Daumke (MDC) John Chodera (MSKCC NY) Gianni de Fabritiis (Barcelona)



